

### III.C.2.a. Individual components of the ORR resolution function

#### 1. Electron burst

The electron burst from ORELA is closely approximated by a square function in time:

$$I_1(t) = \begin{cases} 1/p & \text{for } 0 < t < p \\ 0 & \text{otherwise} \end{cases}, \quad (\text{III C2 a.1})$$

where  $p$  is the burst width in nanoseconds. The value  $1/p$  for this function gives a normalization of unity.

#### 2. Neutron sources

Experiments at ORELA can utilize neutrons from a tantalum target, or collimate on the water moderator surrounding the tantalum target. The two are modeled separately; experiments are assumed to use predominantly one or the other, rather than a mixture of both. (The implications to this representation if neutrons are emitted from both the tantalum target and the water moderator will be studied as time and funding permit.)

##### *ORELA water moderator*

The distribution of flight-path lengths of neutrons in the ORELA water moderator that has the most physical basis [SC92] is the chi-square distribution with  $2(m+1)$  degrees of freedom, where  $m$  is an integer. (Originally  $m = 4$  or  $5$ , but any integer from 1 to 10 will now work; the default value is  $m = 4$ .)

$$I'_{2a}(l) = \frac{l^m}{m! \Lambda^{m+1}} \exp\left(-\frac{l}{\Lambda}\right), \quad (\text{III C2 a.2})$$

where  $l$  is the flight-path-length variable and where the moderation distance  $\Lambda$  is a mean free path. The value of  $\Lambda$  varies with energy as  $\Lambda = \Lambda_0 + \Lambda_1 \ln(E) + \Lambda_2 (\ln(E))^2$  [CC83]. Default values for the  $\Lambda_i$  are chosen to give agreement with [SC92].

### *Tantalum target*

The distribution of flight-path lengths for the tantalum target is somewhat more complicated: Monte Carlo calculations of F. G. Perey [FP92] have shown that this distribution function has the form

$$I'_{2b}(l) = \begin{cases} 0 & \text{for } l < x'_1 \\ u'(l) + v'(l) & \text{for } x'_1 < l < x'_2 \\ u'(l) + w'(l) & \text{for } x'_2 < l < x'_3 \\ w'(l) & \text{for } x'_3 < l \end{cases} \quad (\text{III C2 a.3})$$

where the functions  $u'(l)$ ,  $v'(l)$ , and  $w'(l)$  are defined as

$$u'(l) = N' \exp\{-\varepsilon'^2 (l - x'_0)^2\} \quad \text{for } x'_1 < l < x'_3, \quad (\text{III C2 a.4})$$

$$v'(l) = \frac{N' \alpha' l}{x'_2} \exp\{-\beta(l - x'_2)\} \quad \text{for } x'_1 < l < x'_2, \quad (\text{III C2 a.5})$$

and

$$w'(l) = N' \alpha' \exp\{-\beta'(l - x'_2)\} \quad \text{for } x'_2 < l. \quad (\text{III C2 a.6})$$

The factor  $N'$  is a normalization, and  $\alpha'$  determines the relative strengths of the pieces of the function  $I'_{2b}$ . The relative values of the  $x$ 's are given by

$$x'_1 \leq x'_0 \leq x'_2 \leq x'_3. \quad (\text{III C2 a.7})$$

### **3. Time-of-flight channel width**

Like the electron burst, the time-of-flight channel width may be modeled as a rectangular distribution of width  $c$ . This functional form ignores “electronic” time jitter, which is a reasonable approximation. The time distribution due to the finite channel width is therefore assumed to be

$$I_3(t) = \begin{cases} 1/c & \text{for } 0 < t < c \\ 0 & \text{otherwise} \end{cases}, \quad (\text{III C2 a.8})$$

where the channel width  $c$  may be different for different energy regions.

#### 4. Detectors

Two types of neutron detectors are commonly used for transmission measurements at ORELA: the NE110 (a recoil proton detector) and the lithium glass (for which moderation of neutrons in the detector plays an important role).

##### *NE110 detector*

If  $\delta$  is the thickness of the NE110 detector, then the resolution function appropriate for that detector is

$$I'_{4a}(l) = \begin{cases} \Delta \exp(-\lambda \sigma l) & \text{for } 0 < l < \delta \\ 0 & \text{otherwise} \end{cases} \quad (\text{III C2 a.9})$$

where  $\lambda$  is the number of molecules per volume (molecules per millimeter barn) of the detector (0.0047 for NE110), and  $\sigma(E)$  is the total cross section of the detector material ( $CH_{1.104}$  for NE110). The normalization factor  $\Delta$  is found by setting the integral over all space of  $I'_{4a}$  equal to 1.

##### *Lithium glass detector*

For the lithium glass detector, the resolution function is assumed to have the form

$$I_{4b}(l) = \begin{cases} Dg & \text{for } 0 < t < d \\ D \exp(-f(t-d)) & \text{for } d < t \end{cases}, \quad (\text{III C2 a.10})$$

where  $g$ ,  $f$ , and  $d$  are constants, and  $D$  is chosen to give an integral of unity.